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Effective Two-Group Gamma-Ray Spectra for Thermal Neutron Capture and Prompt Fission



ATOMIC ENERGY COMMISSION CONTRACT W-7405-ENG. 36

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Printed in the United States of America. Available from Clearinghouse for Federal Scientific and Technical Information National Bureau of Standards, U. S. Department of Commerce Springfield, Virginia 22151

Price: Printed Copy \$3.00; Microfiche \$0.65

LA-3619 UC-80, REACTOR TECHNOLOGY TID-4500

of the University of California LOS ALAMOS • NEW MEXICO

Report written: October 1966

Report distributed: February 10, 1967

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by

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ABSTRACT

Effective two-group gamma-ray spectra were determined for thermal neutron capture in sodium, nickel, type 304 stainless steel, and tantalum, as well as for \$235\$U prompt fission gamma rays. The absorbed dose due to uncollided and built-up fluxes in several materials was calculated for varying thicknesses of several intervening shielding materials. A seven-group compilation of capture gamma rays and an analytical fit to the prompt fission gamma-ray spectrum were used as the basis for this study. The resulting function for each combination of shielding and absorbing material was reduced to two exponential functions for areal density ranges with upper limits as great as 560 g/cm². These effective spectra reproduce, to within an average absolute deviation of less than 7.4 percent, the absorbed doses (uncollided and built-up) calculated by the detailed spectra, within the ranges of areal density considered.

ACKNOWLEDGEMENTS

The author is grateful to N. K. Madsen, who did the preliminary FORTRAN IV coding of the calculations and plotting, and to M. E. Battat for his assistance and helpful consultation on this entire work, especially for the determination of total and energy absorption mass attenuation coefficients for tantalum.

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I. INTRODUCTION AND SUMMARY

Basic data on the spectral distribution of gamma rays emitted in neutron radiative capture reactions are usually too detailed to be directly useful for shielding calculations. It thus becomes necessary to collate the data into some convenient energy group structure. Many authors (1-3) have condensed the data into energy group structures of manageable size (five to ten groups) for routine shielding calculations, without compromising the reliability of information extracted from the calculations (e.g., absorbed dose in various materials, including tissue). Mittelman and Liedtke (1) used a five-group structure and a maximum photon energy listing, while Deloume (2) used a ten-group structure. More recently, Troubetzkoy and Goldstein (3) have compiled a seven-group table for most elements of interest; included in the compilation are decay gamma rays from product nuclei with half-lives of the order of hours or less. This seven-group tabulation was used for most of the detailed-spectrum calculations reported herein. It also presents tabulations of line spectra and graphs of differential spectra, which update the work of Bartholomew and Higgs. (4)

Shure and Strobel $^{(5)}$ further reduced the compilations of Deloume $^{(2)}$ and Lanning (unpublished) to effective two-group capture spectra (a three-group spectrum in the case of prompt fission gamma rays) for several common materials of interest to the naval reactors program: iron, zirconium, hafnium, 238 U, and stainless steel (SS). In unpublished work quoted by Guilinger, et al., $^{(6)}$ the author and Schmucker performed the same reduction for Inconel-X (73 w/o nickel, 17 w/o chromium, 10 w/o iron).

This report presents an extension of these results to materials of interest for fast reactors: sodium, nickel, type 304 stainless steel (new spectrum), and tantalum. In addition, a two-group effective 235U prompt fission gamma-ray spectrum has been calculated. A FORTRAN IV program has been written to perform all the calculations, plotting, and curve fitting, so an extension to other spectra is facilitated.

All these effective spectra differ from the many-group spectra in that they were not arrived at by an energy grouping process, but rather by an analytical fitting of typical absorbed dose calculations for various attenuating and absorbing materials. The resulting spectra are then "effective" only in the context of shielding calculations using these materials.

The procedure is, briefly, to calculate for the many-group source spectra the absorbed dose in various materials following varying thicknesses of attenuating material. Both absorbing and attenuating materials cover a wide range of atomic numbers. For each absorbing material (a parameter), the absorbed dose from a unit plane monodirectional source with the many-group spectrum is determined as a function of areal density (g/cm²) of each attenuating material. Each of the resulting functions is then fitted asymptotically by a single exponential, and a corresponding effective energy is determined from the exponent. From the results of many such asymptotic fits, a single energy representative of the attenuation coefficient is chosen, and an attenuation curve for this energy is normalized in the asymptotic region of each function. Residuals are then determined by subtraction of the normalized exponential attenuation curves from the original attenuation curves. Successive representative effective energies are chosen by repeating the fitting process on the residuals. From the intercepts of the exponential curve for each representative energy, effective gamma-ray energy yields are determined. A summary of these yields is presented in Table I.

Most of the capture gamma-ray data accumulated to date have been for thermal neutron capture. These data are normally regarded as being sufficiently accurate for shielding calculations, when used for radiative capture reactions at all energies for which the cross section is significant (up to roughly 100 keV). In some cases, as in the shielding of some fast reactors, these reactions may occur predominantly in the 1 to 100 keV range, opening this assumption to some question. However, for most shield materials of concern, the binding energy of the incident neutron is of the order of 8 MeV, so the total photon energy emitted (essentially equal to the neutron binding energy plus the neutron kinetic energy) does not differ much from the binding energy. This is not to say, however, that the spectrum does not differ significantly. This matter is discussed by Troubetzkoy and Goldstein. (3,7) For lack of better data, the thermal neutron capture spectra are used in practice and have thus far presented no unacceptable discrepancies. Similarly, the prompt gammaray spectrum from thermal neutron-induced fission is used for higher energy fission.

Test calculations of uncollided absorbed doses using the two-group effective spectra show remarkable agreement with comparable many-group calculations over a large range of areal densities. Specific cases are included in the data given in Tables II through VI. These tables list the maximum areal densities over which comparisons were made with manygroup uncollided absorbed doses, the total number of points (i.e., shield thicknesses) at which comparisons were made, the average absolute deviation for the points at which the comparisons were made, the maximum deviation in this range, and the areal density at which the maximum deviation occurs. All these results are for uncollided absorbed dose. Similar results have been obtained with the built-up absorbed dose, as shown in Tables VII through XI. These results were based upon Taylor's buildup parameters, as given by Guilinger, et al. (6) Figure 1 shows a typical attenuation plot (asterisks) in iron, the two effective spectrum curves, and the sum of the two effective spectrum curves. All these curves are for the uncollided absorbed dose in carbon.

Comparisons were also made between the present effective spectra and Shure and Strobel's effective spectra for type 304 stainless steel and prompt fission. The results of uncollided and built-up absorbed dose calculations using the different sets of effective spectra agree reasonably well and indicate that the present effective spectra may be applicable at deeper penetrations than the spectra of Shure and Strobel.

Considering the uncertainties in the raw data from which the many-group compilations are derived, (3) the effective two-group spectra should prove adequate for practical shielding calculations. They should be particularly useful where calculations are performed without the aid of a computer code.

II. CALCULATIONAL PROCEDURE

A. Determination of Many-Group Spectra

The many-group (five or six) capture gamma-ray yields (MeV/capture) were determined by use of the energy groupings of Troubetzkoy and Goldstein. The only modification in these energy groupings was that the >9 MeV group was combined with that for 7-9 MeV, resulting in six groups: 0-1, 1-2, 2-3, 3-5, 5-7, and 7-9+ MeV. Troubetzkoy and Goldstein give the number of photons in each energy group per 100 captures. In cases where the product nucleus has a half-life of the order of hours or less, the decay gamma rays were included in the tabulation. In order to determine the total energy emitted in each group, the discrete gamma-ray energies and intensities were used, wherever available, to find a weighted average energy in each group. Photons in a group which were not accounted for in the line spectra tabulation were then assigned the median energy of the group in the weighting process. In other words, the weighted average energy in the \underline{j} th group, $\widetilde{E}_{\underline{i}}$, is given by

$$\widetilde{E}_{j} = \left\{ \sum_{\nu} E_{j\nu} f_{j\nu} + \left[f_{j} - \sum_{\nu} f_{j\nu} \right] E_{j} \right\} f_{j}^{-1} \qquad (MeV), \qquad (1)$$

where

 $E_{j\nu} \stackrel{\Delta}{\underset{\Lambda}{\longrightarrow}}$ the energy, in MeV, of the $\underline{\nu}$ th line spectrum in group j,

 $\vec{E}_{j} \stackrel{\triangle}{=}$ the median energy, in MeV, of group j,

 $f_{j\nu} \stackrel{\triangle}{=}$ the photon yield, in photons/capture, at the energy of the \underline{v} th line spectrum in group j, and

 $f_{j} \stackrel{\triangle}{=}$ the total photon yield, in photons/capture, of group j.

The gamma-ray energy yield in group j, $Y(E_i)$, is then given by

$$Y(E_j) = f_j \tilde{E}_j$$
 (MeV/capture), (2)

where E_j is the upper limit of energy group j. Values for $E_{j\nu}$ and $f_{j\nu}$ were determined from Table 2 of Troubetzkoy and Goldstein, and values for f_j were determined from their Table 1. In some cases, insufficient line spectra were given (none for tantalum) to enable the use of this weighting scheme, so \widetilde{E}_j was assigned the group median energy, or a value was estimated from Troubetzkoy and Goldstein's curves of the differential capture spectra. The groups are numbered 1 through 6, beginning at the lowest energy (0-1 MeV is group 1, 1-2 MeV is group 2, etc.).

For tantalum capture gamma rays, f_6 = 0, and the differential spectrum curve shows no photon yield above 6 MeV. Thus, the fifth group for tantalum was changed to 5-6 MeV.

To determine the six-group spectrum of type 304 stainless steel, the individual elemental components were weighted by their 2200 m/sec macroscopic capture cross sections as follows:

$$Y(E_{j}) = \frac{\sum_{i}^{\infty} (w/o)_{i} \frac{\sigma_{i}}{A_{i}} Y_{i}(E_{j})}{\sum_{i}^{\infty} (w/o)_{i} \frac{\sigma_{i}}{A_{i}}}$$
 (MeV/capture). (3)

Here

 $(w/o)_{\underline{i}} \stackrel{\Delta}{=}$ the weight percent of the <u>i</u>th element in type 304 stainless steel,

 σ_i the 2200 m/sec microscopic capture cross section, in barns, of the <u>i</u>th material in stainless steel, and

 $A_{i} \stackrel{\Delta}{=}$ the atomic weight of the <u>i</u>th material in stainless steel.

Stainless steel was assumed to have the nominal composition for type 304 stainless steel: (8) 71 w/o iron, 19 w/o chromium, and 10 w/o nickel.

Trace elements such as magnesium and cobalt have lower weight percents

than the allowable variation in weight percent of the three major constituents; they were, therefore, ignored. Values of σ_1 were determined from the compilation of Hughes and Schwartz. (9)

The prompt fission gamma-ray yields for 235U were determined from the analytical fits of the differential prompt fission spectrum data of Maienschein, et al., (10) as given by Goldstein. The analytical expressions for the differential energy spectrum are

I(E) = 26.8 E exp(-2.30E)
$$0.3 \le E \le 1.0$$
 (fission⁻¹), (4)
I(E) = 8.0 E exp(-1.10E) $1.0 \le E \le 7$ (fission⁻¹), (5)

$$I(E) = 8.0 \text{ E exp}(-1.10\dot{E})$$
 $1.0 \le E \le 7$ (fission⁻¹), (5)

where

 $I(E) \stackrel{\Delta}{*}$ the differential prompt fission gamma-ray energy spectrum (MeV/fission-MeV).

These data are for gamma rays emitted with $\approx 5 \times 10^{-8}$ sec after thermal fission of 235 U. Kirkbride $^{(11)}$ has shown that there is no observable difference in the prompt gamma spectral shapes for 235 U, 233 U, and 239 Pu fission. The seven-group spectrum for prompt fission was then determined by integration of Eqs. (4) and (5) over the appropriate ranges; e.g.,

$$Y(1 \text{ MeV}) = 26.8 \int_{0.3}^{1.0} E \exp(-2.30E) dE$$
 (MeV/fission). (6)

In general,

$$Y(E_j) = 8.0 \int_{E_{j-1}}^{E_j} E \exp(-1.10E) dE$$
 (MeV/fission), (7)

for j > 1. Here, we assign the energy released within a group to the upper limit of the group, which is $\mathbf{E_{i}}$. The resulting seven groups were 0.3-1, 1-2, 2-3, 3-4, 4-5, 5-6, and 6-7 MeV.

Both capture and prompt fission gamma-ray yields in any given group were associated with the upper limit of the energy range of that group. Thus, the gamma-ray attenuation and energy absorption coefficients for these groups are characteristic of this upper limit. As was shown in test calculations by Shure and Strobel (5) for the cases they studied, this assumption is usually conservative. In their study, they used the physical properties (attenuation and absorption coefficients) of the median of the group for several test cases. In an indirect comparison with the results using the upper limit of the group, they showed that no significant changes occurred.

Table XII presents the resulting many-group spectra for capture and prompt fission gamma rays, as determined by the procedure described above. Observe that the total energy includes decay gamma rays for product nuclei with half-lives of less than a day. The binding energies are listed for informational purposes from the tabulation of Deloume, (2) who determined a weighted-average value from the following expression:

B.E. =
$$\frac{\sum_{i}^{p_{i} \sigma_{i}} (B.E.)_{i}}{\sum_{i}^{p_{i} \sigma_{i}}}$$
 (MeV), (8)

where

(B.E.) $\stackrel{\Delta}{i}$ the binding energy, in MeV, of the extra neutron in the $\underbrace{(i+1)}_{th}$ isotope,

 P_i the percent abundance of the <u>i</u>th isotope, and σ_i the thermal neutron capture cross section, in barns, of the ith isotope.

No attempt was made to manipulate the many-group spectrum yields to normalize to the binding energies.

B. Reduction to Two-Group Effective Spectra

If the many-group spectrum is known, the absorbed dose can then be determined for a plane monodirectional beam having this spectrum. Consider a semi-infinite slab of material i, with its surface normal to the beam and an infinitesimal absorbing sphere of material k at an areal density of t_1 g/cm² along a normal to the surface. Assuming infinite linear electron stopping power (-dE/dx) for both absorber and shield, the uncollided absorbed dose rate in material k from a unit (one event/cm²) monodirectional beam is given by

$$H_u^k(t_i) = \sum_j \chi_a^k(E_j) Y(E_j) \exp[-\chi^i(E_j) t_i]$$
 (MeV/g), (9)

where

 $\chi_a^k(E_j) \stackrel{\Delta}{=}$ the energy absorption mass attenuation coefficient, in cm²/g, of material k, and

 $\chi^{i}(E_{j}) \stackrel{\triangle}{=} the total mass attenuation coefficient of material i, in cm²/g.$

With the exception of tantalum, values of χ_a^k and χ^i were determined from calculations by Blizard, et al., (12) who used Grodstein's data. (13) For tantalum, these values were determined by M. E. Battat of LASL from the data for pair production and photoelectric effect reported by Storm, et al., (14) and a computed Compton scatter cross section.

For each many-group spectrum, values of $H_u^k(t_i)$ were determined with 20 combinations of attenuating and absorbing materials. There are five of each type of material as follows:

Attenuating Material	Absorbing Material
iron	iron
carbon	carbon
sodium	tantalum
tantalum	concrete
concrete	tissue

The combinations chosen can be seen in the first and third columns of Tables II through XI. Certain combinations were considered to be of marginal practical interest and were, therefore, ignored (e.g., attenuation in tantalum and absorption in tissue). This decision was made to avoid weighting the effective spectrum results by cases of limited concern, because each combination has equal weight in the averaging of effective yields for any given absorbing material. However, once the calculations were coded, it required little extra effort to test the effective spectra on the combinations of limited interest. Such tests give some indication of the validity of these effective spectra for combinations not among those used in the weighting process. Tables XIII and XIV present a summary of these results, for uncollided and built-up absorbed dose, respectively. These results support extending the effective spectra to cases not used in the weighting process. The maximum deviations are generally larger than for the cases used in the weighting process, but are still acceptable.

Consider now any of the five spectra which were reduced. The calculational procedure is as follows: After computing values of $H_u^k(t_i)$ by Eq. (9) for t, over a range varying from 0 up to 560 g/cm² of material i, we fit the resulting values by an exponential function in the asymptotic region. Figure 2 shows a typical asymptotic fit, for the case of a sodium capture spectrum, iron shield, and carbon absorber. The values of $H_{i}^{K}(t_{i})$ are plotted with an asterisk. Letting the asymptotic exponential be given by $A_i^k \exp(-\alpha_i^k t_i)$, we have 20 sets (A_i^k, α_i^k) , corresponding to the 20 combinations of attenuating and absorbing materials. The determination of the asymptotic region was arbitrary and varied depending upon the maximum areal densities considered for each attenuating material. Again, the maximum areal densities are arbitrary parameters, so they were chosen to exceed any shield thickness anticipated for the material in question in a reactor application. An additional restraint on maximum areal densities was a desire to keep reasonable accuracy of the two-group representation in the ranges of most usual interest. The values chosen are listed in Tables II through VI.

All exponential fitting was performed by a general least-squares subroutine with a weighting option, (15) rather than a linear least-squares fit on transformed data. For the asymptotic exponential, $A_{\bf i}^k \exp(-\alpha_{\bf i}^k t_{\bf i}), \text{ the weights were the reciprocal of the square of the ordinates, i.e., } w_1^k(t_{\bf i}) = [H_{\bf u}^k(t_{\bf i})]^{-2}.$ This weighting scheme biases the fit toward the asymptotic region. It is equivalent to minimizing, with a weight of $[H_{\bf u}^k(t_{\bf i})]^{-1}$, percent error. Both parameters, $A_{\bf i}^k$ and $\alpha_{\bf i}^k$, were allowed to vary, and an average value of $\alpha_{\bf i}^k$ was calculated for each material i by the expression

$$\left\langle \alpha_{1}^{k}\right\rangle _{k}=\frac{\sum_{k=1}^{N(1)}\alpha_{1}^{k}}{N(1)},\tag{10}$$

where N depends upon i (see Tables II through VI). From these values of $\left<lpha_i^k
ight>_k$, a representative energy \overline{E}_1 was determined for the attenuation in the asymptotic region, giving a consistent set of values, say $\boldsymbol{\hat{\alpha}}_i$, for this energy. The energy \overline{E}_1 is the higher energy given in Table I for each spectrum studied. Observe that \overline{E}_1 is not in all cases unique for a given attenuating material because the $\chi^{i}(E)$ have minima, but one energy approximately satisfies all attenuating materials considered. The symbol \overline{E}_{i} is used to differentiate the effective energies from the many-group energies E_i . With these values of α_i held fixed, the least-squares fit was repeated with the A, left as free parameters. The same weights as in the previous fit were used. With all asymptotic slopes in a semi-log plot normalized to a particular energy (the $\hat{\alpha}_{i}$ are the total mass attenuation coefficients for the effective energy of the asymptotic exponential), the result of this second fitting process is to determine the intercepts $ilde{\mathtt{A}}_{i}^{k}$ for a fixed slope $\hat{\mathtt{a}}_{i}$. Referring to Eq. (9), observe that the $ilde{\mathtt{A}}_{i}^{k}$ may then be averaged over i to determine an average gamma-ray energy yield at \overline{E}_1 MeV, for any fixed k. That is,

$$Y^{k}(\overline{E}_{1}) = \frac{\sum_{i=1}^{M(k)} \mathring{A}_{i}^{k}}{N(k) \chi_{a}^{k}(\overline{E}_{1})}, \qquad (11)$$

where M depends upon k (see Tables II through VI). By continuing the averaging process to k, the effective yield at energy \overline{E}_1 is found to be

$$\overline{Y}(\overline{E}_1) = \frac{1}{5} \sum_{k=1} Y^k(\overline{E}_1), \qquad (12)$$

where we use the symbol \overline{Y} to differentiate the effective yield from the many-group yields. This averaging process gives equal weight to each absorbing material, not to each combination of attenuating and absorbing materials.

Now, let us define parameters

$$A^{k}(\overline{E}_{1}) = \chi_{a}^{k}(\overline{E}_{1}) \ \overline{Y}(\overline{E}_{1}), \tag{13}$$

which in conjunction with the α_i completely determine a consistent set of asymptotic attenuation curves A^k $\exp(-\alpha_i t_i)$, for each of the 20 combinations (i,k). These curves correspond to the higher-energy yields and energies in Table I. Subtracting these asymptotic attenuation curves from the many-group attenuation curves, we get a set of residual curves. All the residual curves were remarkably close to exponential, accounting in part for the relatively accurate representation of the many-group results by only two effective energies. Fitting each of these residual curves, in turn, by an exponential $B_i^k \exp(-\beta_i^k t_i)$, we get 20 sets (B_i^k, β_i^k) . In this case, the weights chosen were the reciprocal of the ordinates, i.e.,

$$w_2^k(t_i) = [\widetilde{H}_u^k(t_i)]^{-1},$$

where

$$H^{k}(t_{i}) = H_{u}^{k}(t_{i}) - A^{k} \exp(-\hat{a}_{i} t_{i})$$
 (14)

(the residual values). This weighting scheme minimizes percent error. As was true for the asymptotic exponential, both B_i^k and β_i^k were free parameters. Figure 3 shows a typical asymptotic exponential from the consistent set $(A^k,\hat{\alpha}_i)$ for the higher effective energy, the residual values (plotted with a circle), and the curve $B_i^k \exp(-\beta_i^k t_i)$. The case chosen is the same as for Fig. 2. The same process as that for the asymptotic exponential is repeated, and an average value $\langle \beta_i^k \rangle_k$ is determined from an expression analogous to Eq. (10). A representative energy \overline{E}_2 can then in turn be determined. For this energy \overline{E}_2 , we then get a consistent set of values $\hat{\beta}_i$. By constraining the parameter $\hat{\beta}_i$ and using the same weights $w_2^k(t_i)$, we determine the intercepts B_i^k in a manner exactly analogous to the A_i . Figure 4 shows a typical fit of the residual curve, continuing the same case as in Figs. 2 and 3. Again exactly analogous to the asymptotic fit, using equations of the same form as Eqs. (11) through (13), we can determine an effective yield at the second effective energy, $\overline{Y}(\overline{E}_2)$, and an intercept parameter, $B^k(\overline{E}_2)$.

The final two-exponential approximation is shown in Fig. 1 for the typical case we have been following, along with a curve for the sum of the exponentials. Figures 5 through 23 give the same information for all the other cases with a sodium capture spectrum. These curves may be compared with the original calculated points for $H^k_u(t_i)$, which are plotted with an asterisk. All of these results are for the uncollided absorbed dose.

C. Comparison of Two-Group Effective Spectra and Many-Group Spectra

Table XV shows a comparison of the uncollided two-group and sixgroup results for the sample case of a sodium capture spectrum, iron shield, and carbon absorber, where the values given are the ratio

$$R_{u}^{k}(t_{i}) = \frac{H(2-\text{group})}{H(6-\text{group})} = \frac{\sum_{j=1}^{2} \chi_{a}^{k}(\overline{E}_{j}) \overline{Y}(\overline{E}_{j}) \exp[-\chi^{i}(\overline{E}_{j}) t_{i}]}{\sum_{j=1}^{6} \chi_{a}^{k}(E_{j}) Y(E_{j}) \exp[-\chi^{i}(E_{j}) t_{i}]}.$$
 (15)

It is from data such as these that the average absolute deviations and maximum deviations were determined for Tables II through VI. To define these deviations, consider a combination (i,k) and a calculation of R_u^k for N values of areal density $t_{i,\ell}$, as shown in Table XV. The deviation for the ℓ th value of t_i , say $t_{i,\ell}$, is then defined as

$$\varepsilon_{\ell}^{i,k} = R_{u}^{k}(t_{i,\ell}) - 1. \tag{16}$$

The average absolute deviation is then defined by

$$\bar{\varepsilon}^{i,k} = \frac{1}{N} \sum_{\ell=1}^{N} |\varepsilon_{\ell}^{i,k}|, \qquad (17)$$

and the maximum deviation by

$$\varepsilon_{\max}^{i,k} = +\max\left\{\left|\varepsilon_{\ell}^{i,k}\right| : \ell = 1,2,\ldots,N\right\}, \qquad (18)$$

where the sign of $\epsilon_{\max}^{i,k}$ is determined by the sign of the $\epsilon_{\ell}^{i,k}$ that corresponds to the maximum $|\epsilon_{\ell}^{i,k}|$. Likewise, an overall average absolute deviation is defined by

$$\bar{\varepsilon} = 0.05 \sum_{i,k} \bar{\varepsilon}^{i,k}, \qquad (19)$$

and an overall maximum deviation is defined by

$$\varepsilon_{\max} = \pm \max \left\{ \left| \varepsilon_{\max}^{i,k} \right| : \text{all}(i,k) \right\}, \tag{20}$$

where the sign of ϵ_{\max} is determined by the sign of the $\epsilon_{\max}^{i,k}$ that corresponds to the maximum $|\epsilon_{\max}^{i,k}|$.

It is important to observe that $\bar{\epsilon}^{i,k}$ is biased toward the $\epsilon^{i,k}_{\ell}$ for lower values of t_i , because more points were chosen in this region than in the asymptotic region. Table XV is typical of the distribution of the t_i . Also, note that no attempt was made to invoduce a deliberate conservatism into the effective spectra, i.e., to make the effective spectra predict dose rates higher than the many-group spectra. The only bias in the least-squares fit, therefore, was the implicit biasing by the weighting schemes.

As a further comparison of the effective two-group spectra with the many-group spectra, buildup was incorporated in Eq. (15). The form of buildup factor chosen was Taylor's (16) exponential approximation to the moments method data, of the form

$$B(\chi^{i} t_{i}; E) = C_{1}(E) \exp[-c_{1}(E) \chi^{i}(E) t_{i}] + C_{2}(E) \exp[-c_{2}(E) \chi^{i}(E) t_{i}],$$
(21)

where

$$c_2 = 1 - c_1$$
.

The constants C_{ν} and c_{ν} (ν = 1,2) were taken from the tabulation of Guilinger, et al. (6) The equation for the ratio of built-up absorbed doses, corresponding to Eq. (15), is then

$$R^{k}(t_{\mathbf{i}}) = \frac{\sum_{\nu=1}^{2} \sum_{j=1}^{2} C_{\nu}(\overline{E}_{\mathbf{j}}) \chi_{\mathbf{a}}^{k}(\overline{E}_{\mathbf{j}}) \overline{Y}(\overline{E}_{\mathbf{j}}) \exp -[1 + c_{\nu}(\overline{E}_{\mathbf{j}})] \chi^{\mathbf{i}}(\overline{E}_{\mathbf{j}}) t_{\mathbf{i}}}{\sum_{\nu=1}^{2} \sum_{j=1}^{6} C_{\nu}(E_{\mathbf{j}}) \chi_{\mathbf{a}}^{k}(E_{\mathbf{j}}) Y(E_{\mathbf{j}}) \exp -[1 + c_{\nu}(E_{\mathbf{j}})] \chi^{\mathbf{i}}(E_{\mathbf{j}}) t_{\mathbf{i}}}$$
(22)

Table XVI presents a comparison using Eq. (22) for our sample case. In this case, the average absolute deviation was lower with buildup than without, viz., 3.31 and 4.19 percent, respectively. However, the maximum deviation, which occurs at the maximum areal density of 500 g/cm² for both, is slightly higher with buildup, 13.32 vs 12.49 percent. Average absolute deviation and maximum deviation for Eq. (22) are defined exactly as in Eqs. (17) and (18), with the same biasing of $\varepsilon^{-i,k}$. It can be noted from Tables II through XI that the average absolute deviation changes only slightly when buildup is incorporated in the comparisons. In fact, the overall average absolute deviation, $\bar{\epsilon}$, is lower with buildup for nickel and tantalum capture gamma-ray spectra and the prompt fission gamma-ray spectrum. In all cases, incorporating buildup decreased the magnitude of the overall maximum deviation, ε_{max} , for any given source spectrum. The only conclusion that should be drawn from these results is that incorporating buildup did not significantly alter the accuracy of the effective-spectra approximations to the many-group spectra. The changes noted with buildup are within the range one might expect from the errors in the Taylor exponential approximation to the moments data for buildup (p. 73 of Guilinger (6) or pp. 376-378 to Goldstein (7) contain summaries of these errors).

D. Comparison of Previous and Present Effective Spectra for Type 304 Stainless Steel and Prompt Fission

Further comparisons were made between Shure and Strobel's effective spectrum for type 304 stainless steel and the present effective spectrum. These are both two-group effective spectra but were derived using different sources of data for the many-group spectra (see Table XVII), a different assumed composition of type 304 stainless steel, and different combinations of attenuating and absorbing materials. It is thus difficult to separate the causes of discrepancies in calculations using the two different effective spectra. However, as may be noted in Table XVIII, the discrepancies are relatively minor for iron shields, which were used in the

determination of both sets of effective spectra. Table XVIII presents data for the ratio of the present two-group spectrum uncollided dose results to those using Shure and Strobel's effective spectra. Similar results were obtained with built-up absorbed doses. The comparison is good for carbon, sodium, tantalum, and concrete at the lower areal densities. In general, Shure and Strobel used smaller ranges of areal density of attenuating material in determining their effective spectra than did the present study. The new effective spectrum results for type 304 stainless steel would thus be expected to reproduce the manygroup spectrum results better at very deep penetrations. This appears to be the case for carbon, sodium, and concrete, where the present effective spectrum predicts dose rates significantly higher than Shure and Strobel's at the deeper penetrations, yet reproduce the six-group results quite well. The closest comparable result of Shure and Strobel's study is for water which has an attenuation coefficient very similar to concrete in the 7 to 8 MeV range. At the deepest penetration in water which they consider, 300 g/cm², the ratio of tissue dose using their twogroup effective spectrum to that using their ten-group detailed spectrum is 0.791. (5) A similar effect appears in carbon and sodium at penetrations of the order of 300 $\mathrm{g/cm}^2$, indicating that the new effective spectrum may be applicable at deeper penetrations for some of the attenuating materials Shure and Strobel consider.

The comparison of the effective spectra for prompt fission is somewhat more straightforward to analyze, because both spectra are based on the same source spectrum (see Table XVII). As can be seen from Table XVIII, the two effective spectra agree within 11.9 percent over the range of areal density considered. Similar results were obtained for the built-up absorbed doses. This agreement appears to verify the applicability of the three-group effective spectra to attenuating and absorbing materials not used in the determination of the spectra.

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TABLE I
Two-Group Effective Gamma-Ray Energies and Energy Yields

	Ej Gamma-Ray Energy (MeV)	$\overline{Y}(\overline{E}_{j})$ Yield (MeV/capture)
Sodium capture	5.5 2.0	6.09 5.74
Nickel capture	8.0 2.0	8.33 1.62
Stainless steel capture	8.0 2.0	5.86 1.95
Tantalum capture	4.0 1.5	3.76 2.88
		(MeV/fission)
Prompt fission	4.0 1.25	2.31 4.92

TABLE II

Comparison of Effective Spectrum and Six-Group Spectrum
Uncollided Calculations for Sodium Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of <u>Points</u>	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	3.73 4.19 4.05 3.94 4.20	9.86 12.49 9.23 11.41 12.34	460 500 420 500 500
Carbon	300	Iron Carbon Concrete Tissue	24	1.58 1.37 1.33 1.89	- 7.08 4.38 - 4.13 4.52	300 180 300 180
Sodium	260	Iron Carbon Tantalum Concrete Tissue	22	1.43 1.59 1.85 1.44 2.10	3.98 5.51 4.08 4.79 5.53	200 240 180 220 240
Tantalum	200	Iron Tantalum	19 ↓	3.19 3.15	-13.57 -12.29	200 200
Concrete	560	Iron Carbon Concrete Tissue	37	4.74 3.92 4.09 4.27	-22.06 -16.38 -18.58 -16.96	560 560 560 560

Total number of points = 562.

Overall average absolute deviation = 3.14%.

TABLE III Comparison of Effective Spectrum and Six-Group Spectrum Uncollided Calculations for Nickel Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing Material	Number of <u>Points</u>	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	2.31 2.92 2.15 2.22 2.52	5.97 7.49 5.38 6.29 6.84	500 500 500 500 500
Carbon	300	Iron Carbon Concrete Tissue	24	1.86 1.19 1.52 1.25	- 8.16 - 5.22 - 6.94 - 5.93	300 300 300 300
Sodium	300	Iron Carbon Tantalum Concrete Tissue	24	0.71 0.75 1.13 0.50 0.53	- 3.41 1.05 - 4.46 - 2.43 - 1.55	300 50 300 300 300
Tantalum	300	Iron Tantalum	24 •	2.31 1.81	- 8.00 - 6.48	300 300
Concrete	560	Iron Carbon Concrete Tissue	37	1.95 1.90 1.61 1.70	- 8.48 - 4.84 - 6.86 - 5.61	560 560 560 560

Total number of points = 582.

Overall average absolute deviation = 1.73%.

TABLE IV

Comparison of Effective Spectrum and Six-Group Spectrum
Uncollided Calculations for Type 304 Stainless Steel Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of <u>Points</u>	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	1.95 2.19 2.10 1.77 1.94	6.38 7.38 6.04 6.39 6.77	500 500 500 500 500
Carbon	300	Iron Carbon Concrete Tissue	24	1.41 0.69 1.18 0.90	- 6.35 - 3.61 - 5.26 - 4.31	300 300 300 300
Sodium	260	Iron Carbon Tantalum Concrete Tissue	22	0.24 0.58 0.71 0.22 0.29	- 1.13 0.98 - 1.89 - 0.54 0.62	260 220 260 260 20,25
Tantalum	300	Iron Tantalum	24 	2.96 2.43	-12.17 -11.19	300 300
Concrete	560	Iron Carbon Concrete Tissue	37	1.43 1.78 1.25 1.49	- 6.05 3.71 - 4.48 - 3.25	560 280 560 560

Total number of points = 572.

Overall average absolute deviation = 1.46%.

TABLE V

Comparison of Effective Spectrum and Five-Group Spectrum
Uncollided Calculations for Tantalum Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of Points	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	4.26 4.55 4.40 4.46 4.74	10.60 12.00 9.81 11.51 12.25	420 460 400 460 460
Carbon	300	Iron Carbon Concrete Tissue	24	1.85 2.09 2.02 2.49	5.36 6.20 5.89 6.53	200 220 220 220
Sodium	260	Iron Carbon Tantalum Concrete Tissue	22	1.88 1.87 2.37 1.86 2.16	6.14 7.08 5.76 6.75 7.40	240 260 220 260 260
Tantalum ↓	300	Iron Tantalum	24 ↓	5.14 3.80	-12.58 - 9.22	300 280
Concrete	560	Iron Carbon Concrete Tissue	37	4.89 4.25 4.43 4.48	-24.26 -19.34 -21.05 -19.61	560 560 560 560

Total number of points = 572.

Overall average absolute deviation = 3.62%.

TABLE VI

Comparison of Effective Spectrum and Seven-Group Spectrum
Uncollided Calculations for Prompt Fission

Attenuating <u>Material</u>	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of Points	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	400	Iron Carbon Tantalum Concrete Tissue	29	5.48 6.45 4.88 6.08 6.47	13.04 16.47 11.09 15.32 16.92	360 400 320 380 400
Carbon	300	Iron Carbon Concrete Tissue	24	2.97 2.80 2.69 2.70	-16.07 - 9.02 -11.40 - 9.13	300 300 300 300
Sodium	250	Iron Carbon Tantalum Concrete Tissue	24	2.39 3.36 1.87 2.91 3.25	- 4.39 5.92 - 2.74 5.13 6.35	70 235 60 220 235
Tantalum	100	Iron Tantalum	14	4.40 3.66	-14.24 -10.91	100 100
Concrete	340	Iron Carbon Concrete Tissue	26	2.87 3.29 2.98 3.24	-10.79 6.74 - 6.20 7.24	340 220 340 220

Total number of points = 493. Overall average absolute deviation = 3.82%.

TABLE VII Comparison of Effective Spectrum and Six-Group Spectrum Built-Up Calculations for Sodium Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing Material	Number of Points	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	3.33 3.31 4.11 3.22 4.88	13.14 13.32 13.49 13.25 18.42	500 500 500 500 500
Carbon	300	Iron Carbon Concrete Tissue	24	1.50 1.82 1.59 2.32	4.50 5.91 5.27 6.40	240 260 260 260
Sodium	260	Iron Carbon Tantalum Concrete Tissue	22	2.09 1.53 2.94 1.75 1.77	8.33 8.21 8.88 8.26 7.27	260 260 260 260 260
Tantalum	200 ♦	Iron Tantalum	19 ▼	2.64 2.97	-12.70 -12.53	200 200
Concrete	560	Iron Carbon Concrete Tissue	37	4.36 5.13 4.74 4.09	11.48 13.56 12.64 10.95	320 360 340 340

Total number of points = 562.

Overall average absolute deviation = 3.24%.

TABLE VIII Comparison of Effective Spectrum and Six-Group Spectrum Built-Up Calculations for Nickel Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of Points	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	2.52 2.39 2.72 2.13 1.08	6.40 6.70 6.38 6.00 3.75	500 500 500 500 500
Carbon	300	Iron Carbon Concrete Tissue	24	0.93 0.58 0.84 0.86	-3.95 -1.54 -3.06 -2.57	300 300 300 300
Sodium	300	Iron Carbon Tantalum Concrete Tissue	24	0.86 1.47 0.78 0.98 1.61	1.21 2.46 -1.60 1.34 -2.23	200 260 10 240 160
Tantalum	300	Iron Tantalum	24 •	2.94 2.11	6.23 4.53	100 120
Concrete	560	Iron Carbon Concrete Tissue	37	1.85 3.06 2.14 0.88	3.63 5.55 4.14 -2.24	280 320,340 300,320 560

Total number of points = 582. Overall average absolute deviation = 1.71%.

TABLE IX

Comparison of Effective Spectrum and Six-Group Spectrum
Built-Up Calculations for Type 304 Stainless Steel Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of <u>Points</u>	Average Absolute Deviation (%)	Maximum Deviation(%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	1.56 1.43 2.01 1.31 1.11	5.65 5.21 5.97 4.82 2.70	500 500 500 500 500
Carbon	300	Iron Carbon Concrete Tissue	24	0.76 0.25 0.69 0.79	-2.66 -0.60 -1.98 -1.70	300 300 300 300
Sodium	260	Iron Carbon Tantalum Concrete Tissue	22	0.68 0.98 0.87 0.61 1.77	1.32 2.23 -1.83 1.32 -3.27	260 260 10 260 160
Tantalum ↓	300	Iron Tantalum	24 •	3.55 2.53	-7.43 -6.85	300 300
Concrete	560 	Iron Carbon Concrete Tissue	37	2.13 3.25 2.40 1.05	4.36 6.29 4.88 -2.31	320,340 380 360 140

Total number of points = 572.

Overall average maximum deviation = 1.56%.

TABLE X Comparison of Effective Spectrum and Five-Group Spectrum Built-Up Calculations for Tantalum Capture

Attenuating Material	Maximum Areal Density (g/cm ²)	Absorbing Material	Number of Points	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	500	Iron Carbon Tantalum Concrete Tissue	34	3.38 3.44 3.84 3.47 3.62	11.97 12.54 11.49 12.39 11.13	500 500 500 500 500
Carbon	300	Iron Carbon Concrete Tissue	24	2.82 3.10 2.93 1.89	5.53 6.40 6.10 6.50	260 300 280 280
Sodium	260	Iron Carbon Tantalum Concrete Tissue	22	1.51 1.58 2.41 1.44 1.56	4.93 4.53 5.28 4.71 4.76	260 260 260 260 260
Tantalum	300	Iron Tantalum	24 ♦	7.36 6.43	-19.08 -15.90	280 260
Concrete	560	Iron Carbon Concrete Tissue	37	3.10 3.08 2.99 2.93	-11.77 7.41 - 8.89 - 8.28	560 340 560 560

Total number of points = 572.

Overall average absolute deviation = 3.19%.

TABLE XI

Comparison of Effective Spectrum and Seven-Group Spectrum
Built-Up Calculations for Prompt Fission

Attenuating <u>Material</u>	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of <u>Points</u>	Average Absolute Deviation (%)	Maximum Deviation (%)	Areal Density for Max.Dev. (g/cm ²)
Iron	400	Iron Carbon Tantalum Concrete Tissue	29	4.41 5.12 4.13 5.31 4.94	15.10 15.55 14.67 15.55 14.79	400 400 400 400 400
Carbon	300	Iron Carbon Concrete Tissue	24	2.79 3.32 2.56 2.26	- 4.93 - 5.96 4.79 5.20	90 100 260 260
Sodium	250	Iron Carbon Tantalum Concrete Tissue	24	2.64 3.27 2.86 3.13 2.59	- 6.39 - 8.03 - 4.52 - 6.80 - 5.86	115 130 100 130 130
Tantalum	100	Iron Tantalum	14 ↓	3.78 3.10	-14.61 -12.37	100 100
Concrete	340	Iron Carbon Concrete Tissue	26	2.66 3.63 3.49 2.90	- 5.83 7.47 6.74 5.96	120 300 300 300

Total number of points = 493.

Overall average absolute deviation = 3.51%.

TABLE XII Many-Group Capture and Prompt Fission Gamma-Ray Spectra

		Y(E _j) MeV/capture								
	Energy		Capture Material							
Group	Range(a) (MeV)	Sodium(b)	Chromium	Iron	Nickel (c)	Type 304 Stainless Steel	Tantalum (d)			
1	0-1	0.806	0.425	0.375	0.42	0.392	1.03			
2	1-2	1.90	0.615	0.96	0.60	0.832	1.48			
3	2-3	4.63	0.525	0.675	0.575	0.628	1.65			
4	3-5	2.63	0.48	0.867	0.92	0.791	2.20			
5	5-7	1.89	1.39	1.50	2.11	1.56	0.275			
6	7-9+	0.0	3.82	3.11	5.31	3.59				
	Total energy	11.86	7.26	7.49	9.94	7.80	6.64			
	Binding energy	6.96	9.22	7.65	8.87		6.03			

		Y(E _j) MeV/fission
		Prompt Fission
1	0.3-1	2.62
2	1-2	2.28
3	2-3	1.30
4	3-4	0.610
5	4-5	0.263
6	5-6	0.107
7	6-7	0.0423
•	Total energy	7.22

⁽a) The upper energy of the group, E_j , is used to determine the attenuation and absorption coefficient for that group.

(b) Includes ²⁴Na decay gamma rays (4.12 MeV).

(c) Includes ⁶⁵Ni decay gamma rays.

(d) Includes ^{182m}Ta decay gamma rays.

TABLE XIII

Comparison of Effective Spectrum and Many-Group Spectrum
Uncollided Calculations for Five Cases Not Used in the Weighting Process

Source	Attenuating	Maximum Areal Density	Absorbing	Number of	Average Absolute Deviation	Maximum Deviation	Areal Density for Max.Dev.
Spectrum	<u>Material</u>	(g/cm^2)	<u>Material</u>	Points	(%)	(%)	(g/cm^2)
Sodium capture	Carbon Tantalum	300 200	Tantalum Carbon Concrete Tissue	24 19 19 19	2.25 3.43 3.31 3.41	- 9.3 -14.79 -14.14 -13.28	300 200 200 200
.*	Concrete	560	Tantalum	37	5.72	-24.75	560
Nickel capture	Carbon Tantalum Concrete	300 560	Tantalum Carbon Concrete Tissue Tantalum	24 24 24 24 37	2.50 3.15 2.70 3.02 2.31	- 9.39 -10.49 - 9.84 -10.06 -10.01	300 300 300 300 560
304 Stainless steel capture	Carbon Tantalum Concrete	300 	Tantalum Carbon Concrete Tissue Tantalum	24 24 24 24 37	2.20 3.85 3.43 3.74 1.69	- 7.49 -13.68 -13.31 -13.14 - 7.55	300 300 300 300 560
Tantalum capture	Carbon Tantalum Concrete	300 560	Tantalum Carbon Concrete Tissue Tantalum	24 24 24 24 37	2.09 6.44 6.05 6.22 5.88	5.07 -16.19 -14.97 -15.30 -27.95	180 300 300 300 560
Prompt fission	Carbon Tantalum Concrete	300 100 1 340	Tantalum Carbon Concrete Tissue Tantalum	24 14 14 14 26	3.62 5.04 5.02 4.93 3.16	-21.10 -16.11 -14.76 -15.52 -15.78	300 100 100 100 340

TABLE XIV

Comparison of Effective Spectrum and Many-Group Spectrum
Built-Up Calculations for Five Cases Not Used in the Weighting Process

Source Spectrum	Attenuating <u>Material</u>	Maximum Areal Density (g/cm ²)	Absorbing <u>Material</u>	Number of Points	Average Absolute Deviation (%)	Maximum Deviation(%)	Areal Density for Max.Dev. (g/cm ²)
Sodium capture	Carbon	300	Tantalum	24	2.02	4.75	220
	Tantalum	200	Carbon	19	2.90	-12.63	200
į	1	-1-	Concrete	19	2.84	-12.39	200
l l	1	1	Tissue	19	3.49	-11.05	200
†	Concrete	560	Tantalum	37	4.73	11.36	300
Nickel capture	Carbon	300	Tantalum	24	1.52	- 4.95	300
1 -	Tantalum	1	Carbon	24	4.14	8.59	100
1	1		Concrete	24	3.59	7.49	100
	. ↓	+	Tissue	24	4.15	8.73	100
†	Concrete	560	Tantalum	37	1.70	- 3.27	560
304 Stainless	Carbon	300	Tantalum	24	1.38	- 3.50	300
steel capture	Tantalum	ļ	Carbon	24	5.03	9.65	100
[-	İ	l	Concrete	24	4.41	8.42	100
j	. ↓	. ↓	Tissue	24	5.13	9.94	100
. •	Concrete	560	Tantalum	37	1.88	3.78	300
Tantalum capture	Carbon	300	Tantalum	24	3.36	5.24	240
ı	Tantalum	l	Carbon	24	8.45	-22.37	280
İ		ł	Concrete	24	8.13	-21.25	280
		į.	Tissue	24	8.35	-21.51	280
\	Concrete	560	Tantalum	37	3.88	-15.14	560
Prompt fission	Carbon	300	Tantalum	24	3.16	- 4.66	25
1	Tantalum	100	Carbon	14	4.31	-14.79	100
İ	ł		Concrete	14	4.78	-13.65	100
	Į.	. ↓	Tissue	14	4.41	-14.32	100
†	Concrete	340	Tantalum	26	2.69	4.36	260

TABLE XV
Uncollided Absorbed Dose

Spectrum Sodium Capture Shield Material Iron

Absorption Material Carbon

Shield Thickness (g/cm ²)	H(2-group)/H(6-group)
0	0.9971
5	0.9984
10	0.9991
15	0.9992
20	0.9988
25	0.9980
30	0.9968
40	0.9938
50	0.9904
60	0.9868
70	0.9835
80	0.9807
90	0.9784
100	0.9769
120	0.9763
140	0.9787
160	0.9841
180	0.9918
200	1.0015
220	1.0124
240	1.0242
260	1.0364
280	1.0485
300	1.0603
320	1.0715
340	1.0818
360	1.0912
380	1.0995
400	1.1067
420	1.1127
440	1.1175
460	1.1211
480	1.1235
500	1.1249

Number of cases = 34.

Total deviation = 1.4245.

Average deviation = 0.0419.

Maximum deviation = 0.1249.

TABLE XVI

Built-Up Absorbed Dose

Spectrum Sodium Capture

Shield Material Iron

Absorption Material Carbon

Shield Thickness (g/cm ²)	H(2-group)/H(6-group)
0	0.9971
5	0.9932
10	0.9924
15	0.9933
20	0.9949
25	0.9968
30	0.9987
40	1.0015
50	1.0027
60	1.0021
70	1.0000
80	0.9965
90	0.9921
100	0.9872
120	0.9768
140	0.9675
160	0.9607
180	0.9570
200	0.9567
220	0.9598
240	0.9660
260	0.9748
280	0.9858
300	0.9985
320	1.0124
340	1.0271
360	1.0421
380	1.0571
400	1.0718
420	1.0860
440	1.0993
460	1.1117
480	0.1230
500	1.1332

Number of cases = 34.

Total deviation = 1.1243.

Average deviation = 0.0331.

Maximum deviation = 0.1332.

TABLE XVII

Gamma-Ray Energy Yields for Many-Group Spectra
(MeV/capture)

Energy Range	Type 304 Stain1	ess Steel	Prompt Fission			
(MeV)	Shure & Strobel	This Report	Shure & Strobel	This Report		
0-1	1.11	0.392	2.62 ^(a)	2.62		
1-2	0.896	0.832	2.26	2.28		
2-3	0.171	0.628	1.30	1.30		
3-4	0.338)	0.701	0.61	0.610		
4-5	0.486 0.824	0.791	0.26	0.263		
5–6	0.964)	1 56	0.11	0.107		
6–7	0.653 1.617	1.56	0.04	0.0423		
78 89	2.17 1.17 } 3.564	3.59				
9-10	0.224)					

⁽a) Shure and Strobel break this group into four subgroups: 0-0.4, 0.4-0.6, 0.6-0.8, and 0.8-1.0 MeV.

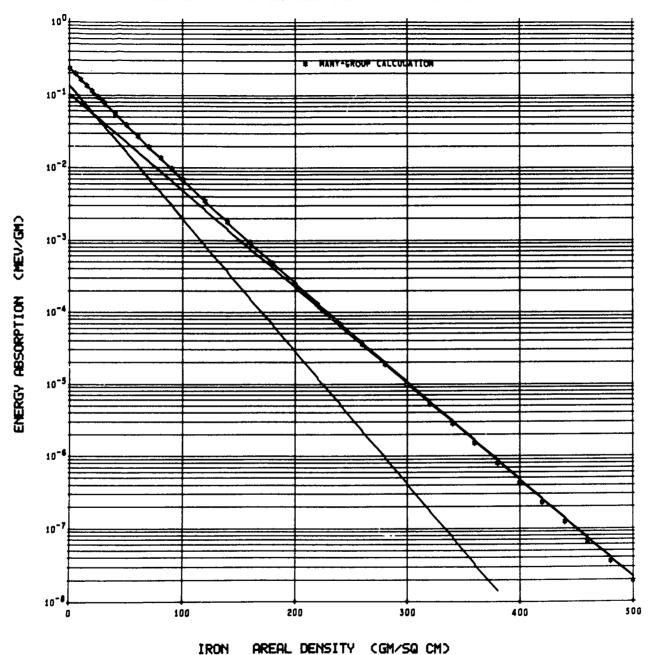
TABLE XVIII Comparison of Uncollided Doses Using Two Different Effective Spectra for 304 Stainless Steel and for Prompt Fission

<u>Spectrum</u>	Attenuating Material	Maximum Areal Density (g/cm ²)	Average Absolute Deviation(a) (%)	Maximum Deviation(b)(%)	Absorbing Material with Max. Dev.	Areal Density for Max.Dev. (g/cm ²)
304 Stainless	Iron	480	2.7	- 9.8	Carbon	0
steel	Carbon	200	5.7	21.7	Iron	200
	Sodium	200	4.8	14.91	Iron	200
	Tantalum	300	3.8	-10.9	Tantalum	300
	Concrete	300	9.5	35.5	Iron	300
Prompt fission	Iron	480	4.6	11.9	Tanta1um	480
	Carbon	300	3.5	11.3	Iron	300
	Sodium	300	3.4	10.7	Tantalum	300
	Tantalum	100	2.6	-10.0	Iron	100
	Concrete	300	3.3	10.5	Iron	300

(a) Average absolute deviation here is defined by

$$\frac{1}{N(i)} \sum_{k=1}^{N(i)} \overline{\epsilon}^{i,k}$$

 $\frac{1}{N(i)}\sum_{k=1}^{N(i)}\bar{\epsilon}^i,k.$ (b) Maximum deviation here is defined by $\pm \max\{|\epsilon_{\max}^{i,k}|: k=1,2,\ldots,N(i)\}.$



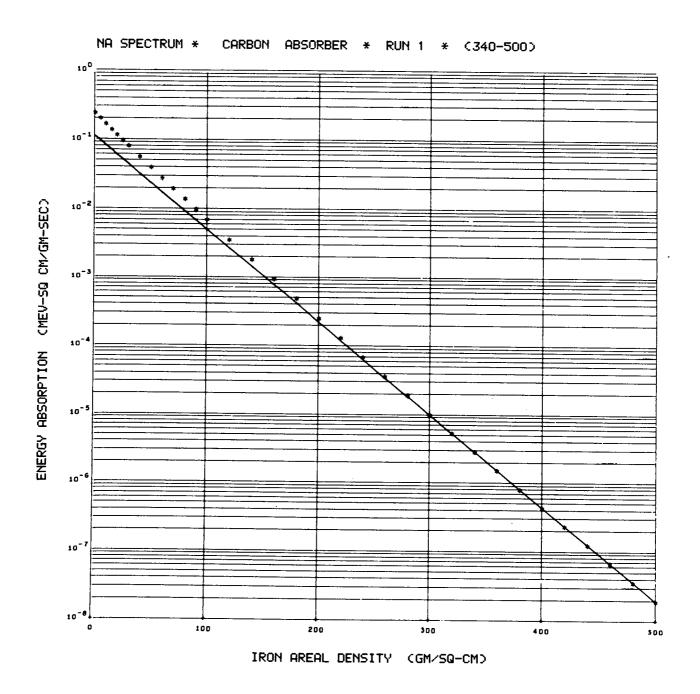


Fig. 2

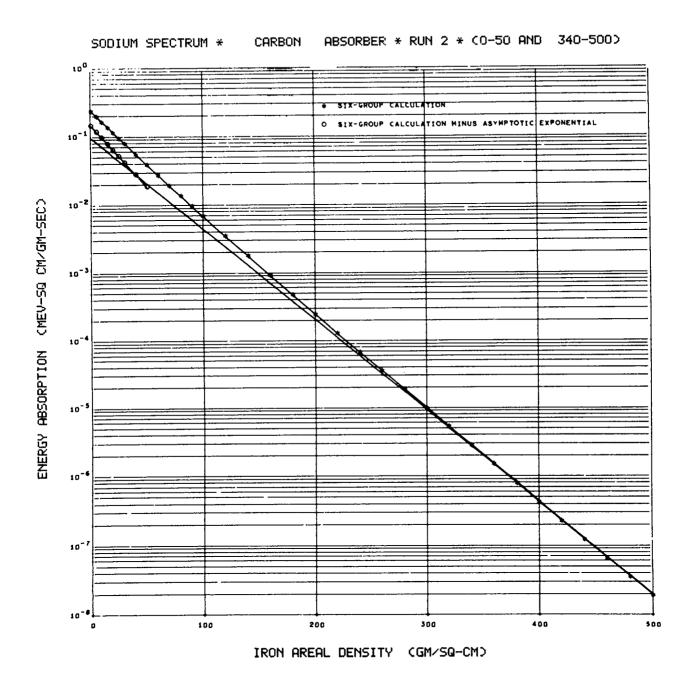


Fig. 3

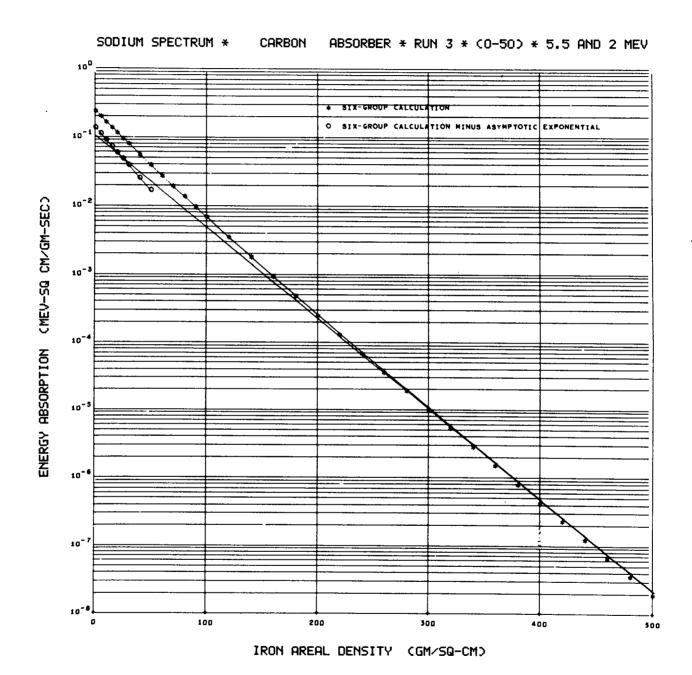
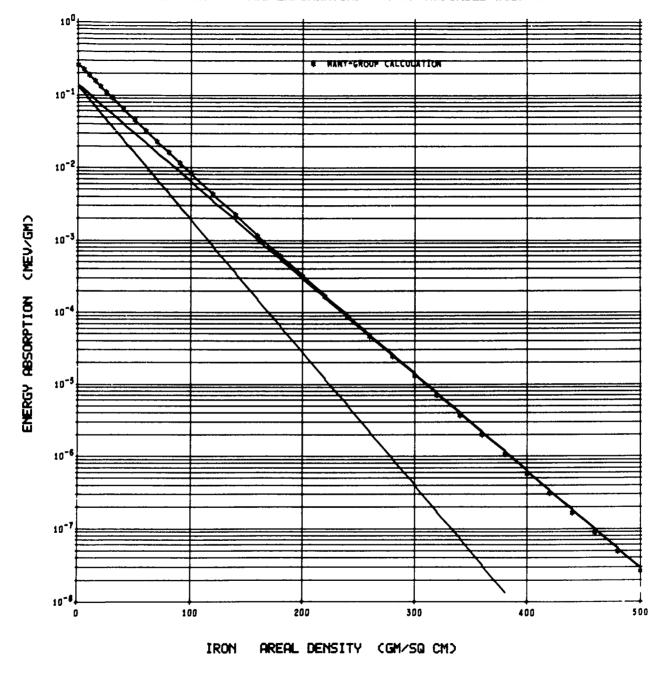
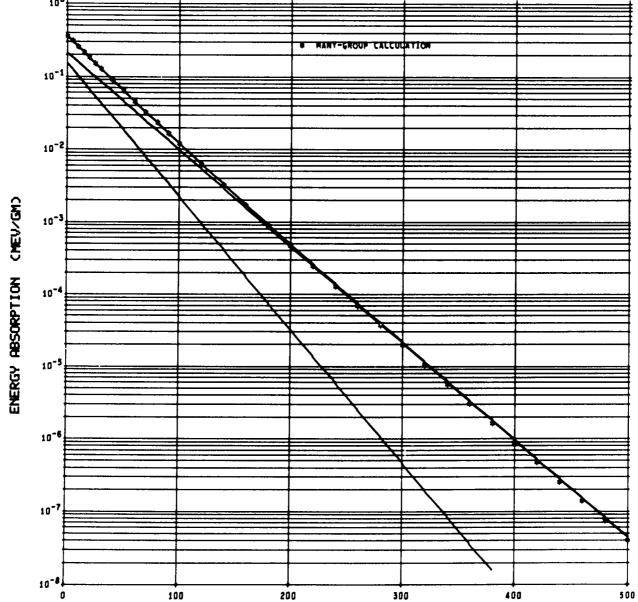


Fig. 4





AREAL DENSITY (GM/SQ CM)

IRON

MARY-GROUP CALCULATION

IRON AREAL DENSITY (GM/SQ CM)

300

400

500

200

Fig. 7

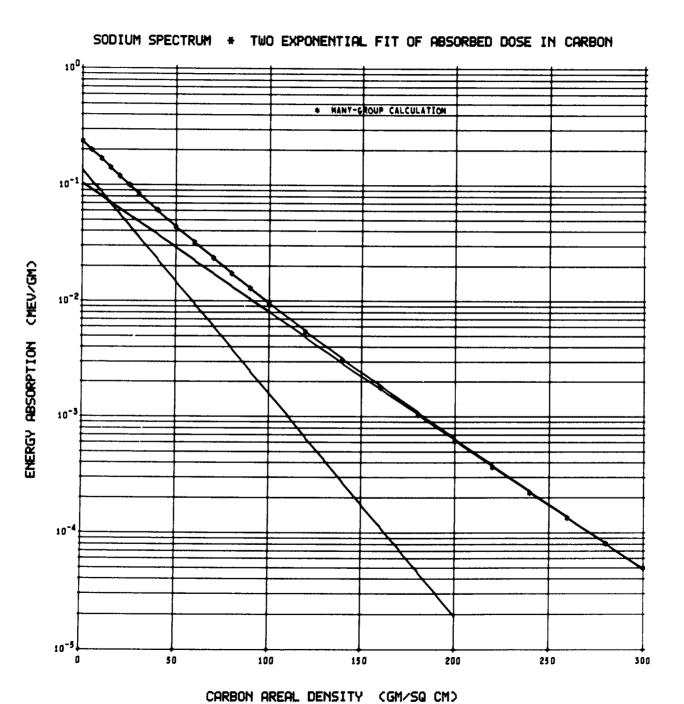
48

100

10-8

AREAL DENSITY (GM/SQ CM)

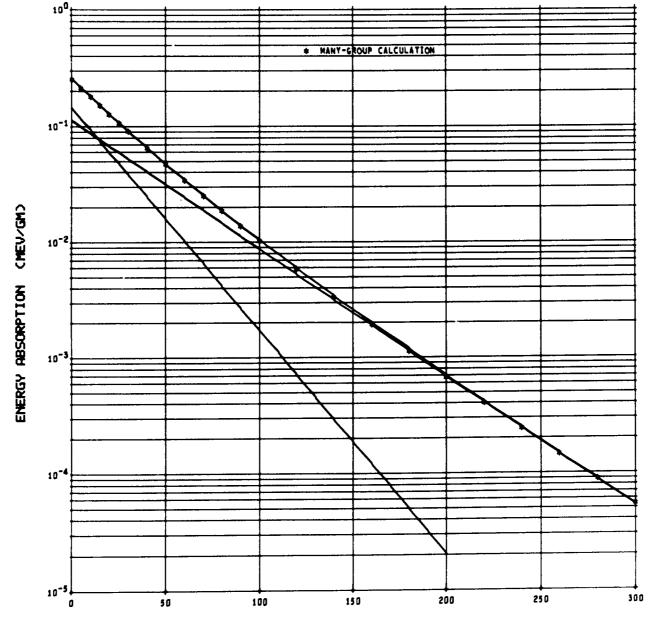
IRON



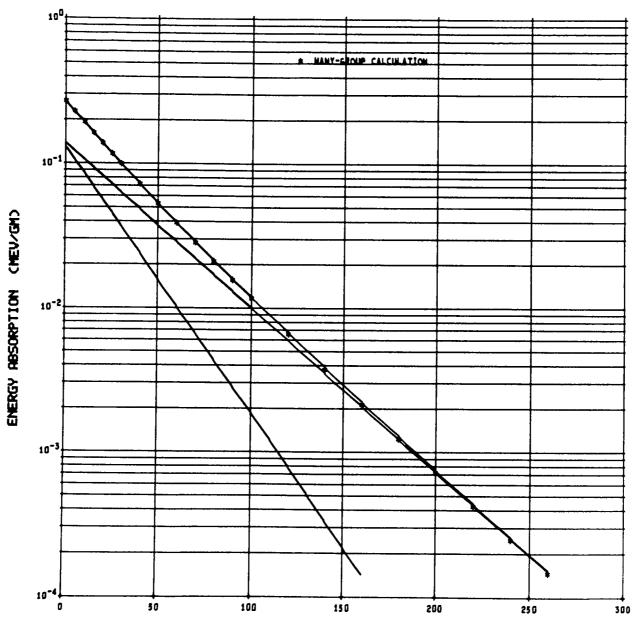
CARBON AREAL DENSITY (GM/SQ CM)

8. 1

52

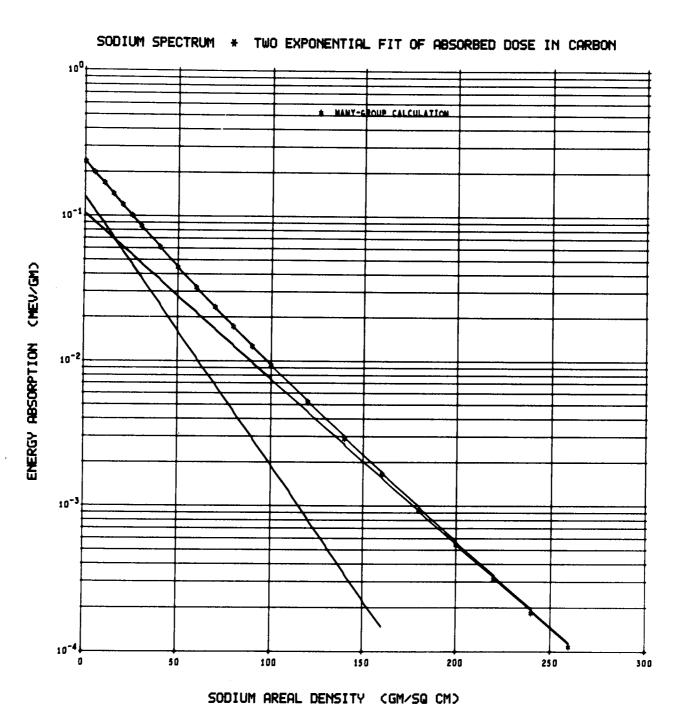


CARBON AREAL DENSITY (GM/SQ CM)

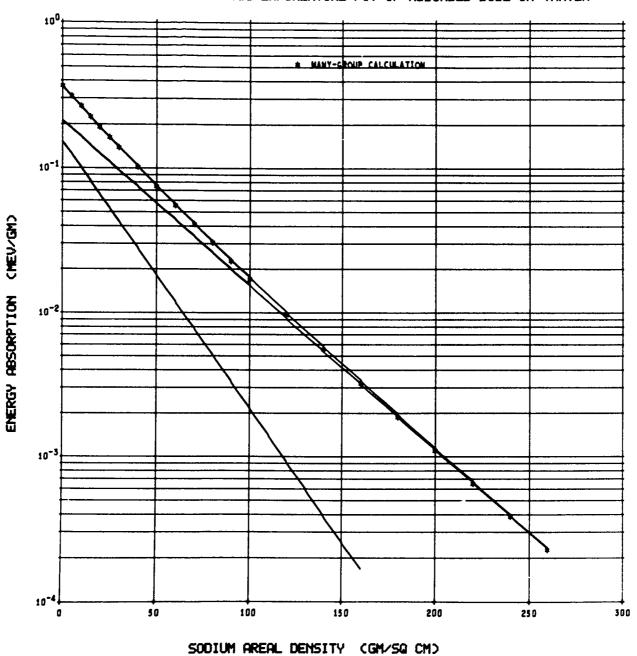


SODIUM AREAL DENSITY (GM/SQ CM)

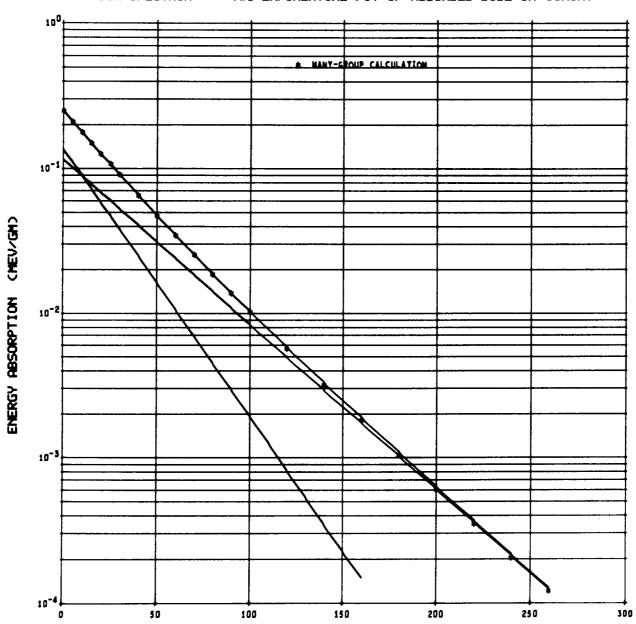












SODIUM AREAL DENSITY (GM/SQ CM)

SODIUM AREAL DENSITY (GM/SQ CM)

80

60

40

20

100

TANTLM AREAL DENSITY (GM/SQ CM)

120

140

160

180

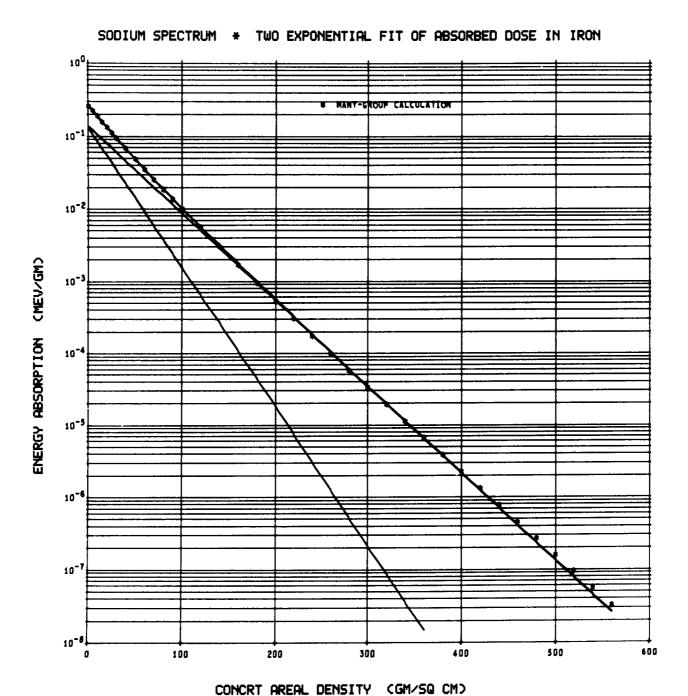
200

ig. 1

TANTLM AREAL DENSITY (GM/SQ CM)

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300

CONCRT AREAL DENSITY (GM/SQ CM)

400

500

600

100

10-8

100

10-8

100

64

CONCRT AREAL DENSITY (GM/SQ CM)

300

400

500

200